

Solutions to Problem Set No. 7

UBC Metro Vancouver Physics Circle 2018-2019

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Problem 1 — Intense Fireworks

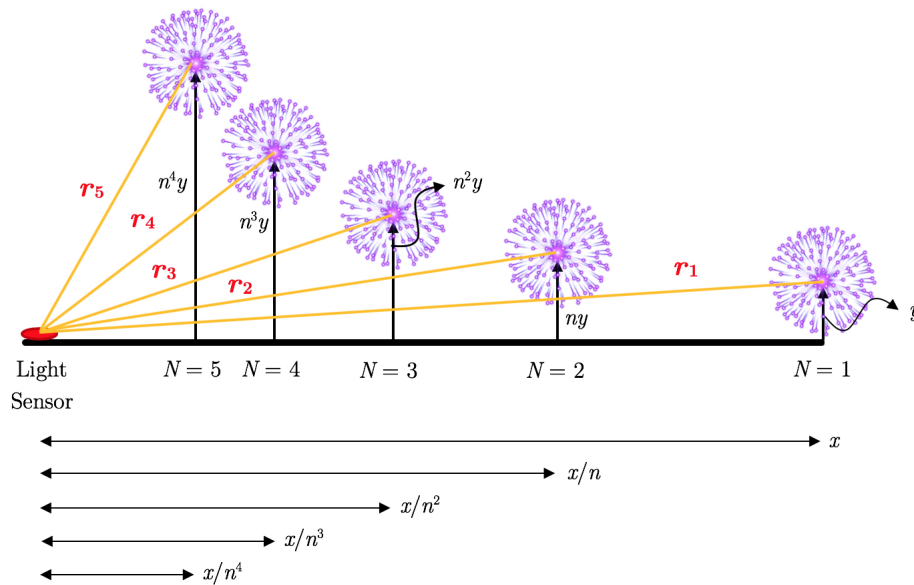
Intensity is the amount of energy a source conveys per unit time across a surface of unit area.

$$I = \frac{P}{A} = \frac{P}{4\pi r^2}$$

Thus, intensity follows an inverse-square law (assuming a source generating constant power),

$$I \propto \frac{1}{r^2}$$

where intensity decreases significantly as the energy is dispersed in space. When we refer to the diagram below,



each firework explosion results in a right angle triangle, where the path of light to the light sensor is the hypotenuse of the triangle. For example, the intensity of the first shot ($N = 1$) can be described as

$$I_1 \propto \frac{1}{(r_1)^2} = \frac{1}{x^2 + y^2}$$

In order to show that $(I_{N+1}/I_N) = 1/n^2$ as N increases indefinitely, we must come up with a general expression of this ratio in terms of N only. An important note to take into account is that when we are taking the ratios, all constants are cancelled out since the only variable changing is r . Let's begin with calculating some r values for different shots.

$$(r_1)^2 = x^2 + y^2$$

$$\begin{aligned} (r_2)^2 &= \left(\frac{x}{n}\right)^2 + (ny)^2 \\ &= \frac{x^2}{n^2} + n^2y^2 \\ &= \frac{x^2 + n^4y^2}{n^2} \end{aligned}$$

$$\begin{aligned} (r_3)^2 &= \left(\frac{x}{n^2}\right)^2 + (n^2y)^2 \\ &= \frac{x^2}{n^4} + n^4y^2 \\ &= \frac{x^2 + n^8y^2}{n^4} \end{aligned}$$

$$\begin{aligned} (r_4)^2 &= \left(\frac{x}{n^3}\right)^2 + (n^3y)^2 \\ &= \frac{x^2}{n^6} + n^6y^2 \\ &= \frac{x^2 + n^{12}y^2}{n^6} \end{aligned}$$

Therefore, based on the pattern that we see above,

$$\begin{aligned}
(r_N)^2 &= \left(\frac{x}{n^{N-1}}\right)^2 + (n^{N-1}y)^2 \\
&= \frac{x^2}{n^{2N-2}} + n^{2N-2}y^2 \\
&= \frac{x^2 + n^{4N-4}y^2}{n^{2N-2}}
\end{aligned}$$

Now that we have the r_N expression for any number of shot N , we can determine our consecutive intensity ratio:

$$\begin{aligned}
\frac{I_{N+1}}{I_N} &= \frac{(r_N)^2}{(r_{N+1})^2} \\
&= \frac{(x^2 + n^{4N-4}y^2) / (n^{2N-2})}{(x^2 + n^{4(N+1)-4}y^2) / (n^{2(N+1)-2})} \\
&= \frac{x^2 + n^{4N-4}y^2}{n^{2N-2}} \cdot \frac{n^{2N}}{x^2 + n^{4N}y^2} \\
&= n^{2N-2N+2} \cdot \frac{x^2 + n^{4N-4}y^2}{x^2 + n^{4N}y^2} \\
&= n^2 \cdot \frac{x^2 + n^{4N-4}y^2}{x^2 + n^{4N}y^2}
\end{aligned}$$

If we apply the limit as N increases indefinitely, we obtain:

$$\begin{aligned}
\lim_{N \rightarrow \infty} \frac{I_{N+1}}{I_N} &= n^2 \lim_{N \rightarrow \infty} \left(\frac{x^2 + n^{4N-4}y^2}{x^2 + n^{4N}y^2} \right) \\
&= n^2 \lim_{N \rightarrow \infty} \left(\frac{n^{4N-4}y^2}{n^{4N}y^2} \right) \\
&= n^2 \lim_{N \rightarrow \infty} (n^{4N-4-4N}) \\
&= n^2 \cdot n^{-4}
\end{aligned}$$

$$\boxed{\lim_{N \rightarrow \infty} \frac{I_{N+1}}{I_N} = \frac{1}{n^2}}$$

Problem 2 — A Sticky Situation

1. In order to find the expression for $\frac{x}{D}$, our goal is to find x (as D can be easily found later). This quantity is the horizontal range of the combined masses from the half-way point. Since this is within the horizontal direction, we can define it to be

$$x \equiv (\text{horizontal velocity of combined masses}) \times (\text{time it takes to hit the ground})$$

$$x = u_x \tau$$

Let's begin with the first unknown, u_x , defined as the horizontal velocity of the combined masses. To find this, we must use conservation of momentum. First, let's find an expression for h then use that to find the velocity of impact of M as it drops a height of h . In the vertical direction:

$$v^2 = v_0^2 + 2ad$$

$$(v_{m,f})^2 = (v_{m,i})^2 - 2gh$$

$$0 = (v \sin \theta)^2 - 2gh$$

$$h = \frac{(v \sin \theta)^2}{2g} \quad (\text{Eqn. 1})$$

$$v^2 = v_0^2 + 2ad$$

$$(v_{M,f})^2 = (v_{M,i})^2 + 2gh$$

$$(v_{M,f})^2 = 2g \left(\frac{(v \sin \theta)^2}{2g} \right) \quad (\text{Using Eqn. 1})$$

$$v_{M,f} = v \sin \theta$$

This essentially means that the final vertical velocity of M is the initial vertical velocity of m , which makes sense as they both traverse the same displacement.

It is now time to use conservation of momentum within the x -direction:

$$\begin{aligned}
 p_i &= p_f \\
 p_{m,i} + \cancel{p_{M,i}}^0 &= p_f \\
 mv \cos \theta &= (m + M)u_x \\
 \frac{mv \cos \theta}{m + M} &= u_x
 \end{aligned}$$

If $\alpha = \frac{m}{M}$, then our horizontal velocity will be:

$$u_x = \frac{\alpha}{\alpha + 1} v \cos \theta \quad (\text{Eqn. 2})$$

Now, let's use conservation of momentum in the y -direction to find the vertical velocity of the combined masses.

$$\begin{aligned}
 p_i &= p_f \\
 \cancel{p_{m,i}}^0 + p_{M,i} &= p_f \\
 Mv \sin \theta &= (m + M)u_y \\
 \frac{Mv \sin \theta}{m + M} &= u_y
 \end{aligned}$$

Again, if $\alpha = \frac{m}{M}$, then our vertical velocity will be:

$$u_y = \frac{v \sin \theta}{\alpha + 1} \quad (\text{Eqn. 3})$$

Let's use this information to calculate τ , the time it takes the combined masses to hit the ground.

$$\begin{aligned}
 d &= v_0 t + \frac{1}{2} a t^2 \\
 h &= u_y \tau + \frac{1}{2} g \tau^2 \\
 0 &= \frac{1}{2} g \tau^2 + \frac{v \sin \theta}{\alpha + 1} \tau - \frac{(v \sin \theta)^2}{2g} \quad (\text{Using Eqns. 1 and 3})
 \end{aligned}$$

To solve for τ , we must use the quadratic equation.

$$\begin{aligned}
\tau &= \frac{-\frac{v \sin \theta}{\alpha + 1} \pm \sqrt{\frac{(v \sin \theta)^2}{(\alpha + 1)^2} - 4 \left(\frac{1}{2}g\right) \left(-\frac{(v \sin \theta)^2}{2g}\right)}}{2 \left(\frac{1}{2}g\right)} \\
&= \frac{-\frac{v \sin \theta}{\alpha + 1} + \sqrt{\frac{(v \sin \theta)^2}{(\alpha + 1)^2} + (v \sin \theta)^2}}{g} && \text{(Choosing only the } \oplus \text{ root.)} \\
&= \frac{v \sin \theta}{g} \left(-\frac{1}{\alpha + 1} + \sqrt{\frac{1}{(\alpha + 1)^2} + 1} \right) && \text{(Factoring } v \sin \theta \text{ out.)} \\
&= \frac{v \sin \theta}{g} \left(-\frac{1}{\alpha + 1} + \sqrt{\frac{1 + (\alpha + 1)^2}{(\alpha + 1)^2}} \right)
\end{aligned}$$

By combining everything together, we achieve:

$$\tau = \frac{v \sin \theta}{g} \left(\frac{\sqrt{(\alpha + 1)^2 + 1} - 1}{\alpha + 1} \right) \quad \text{(Eqn. 4)}$$

Now that we've calculated for both u_x (Eqn. 2) and τ (Eqn. 4), we can simply multiply the two together to determine the horizontal displacement of the combined masses:

$$x = \frac{v^2 \sin \theta \cos \theta}{g} \cdot \frac{\alpha \left(\sqrt{(\alpha + 1)^2 + 1} - 1 \right)}{(\alpha + 1)^2} \quad \text{(Eqn. 5)}$$

In the original diagram, D is half of the range, R . To calculate R :

$$\begin{aligned}
d &= v_0 t + \frac{1}{2} a t^2 && R = (v \cos \theta) t \\
0 &= (v \sin \theta) t - \frac{1}{2} g t^2 && R = \frac{2v^2 \sin \theta \cos \theta}{g} \\
t &= \frac{2v \sin \theta}{g} && \therefore D = \frac{v^2 \sin \theta \cos \theta}{g} \quad \text{(Eqn. 6)}
\end{aligned}$$

Finally, we obtain what we are looking for when we divide Eqn. 5 by Eqn. 6:

$$\boxed{\frac{x}{D} = \frac{\alpha \left(\sqrt{(\alpha + 1)^2 + 1} - 1 \right)}{(\alpha + 1)^2}}$$

2. If $f(\alpha) = \frac{\alpha \left(\sqrt{(\alpha + 1)^2 + 1} - 1 \right)}{(\alpha + 1)^2}$, then:

$$\begin{aligned} \lim_{\alpha \rightarrow \infty} f(\alpha) &= \lim_{\alpha \rightarrow \infty} \frac{\alpha \left(\sqrt{(\alpha + 1)^2 + 1} - 1 \right)}{(\alpha + 1)^2} \\ &= \lim_{\alpha \rightarrow \infty} \frac{\alpha \left(\sqrt{\alpha^2 + 2\alpha + 2} - 1 \right)}{\alpha^2 + 2\alpha + 1} \\ &= \lim_{\alpha \rightarrow \infty} \frac{\alpha \left(\alpha \sqrt{1 + \frac{2}{\alpha} + \frac{2}{\alpha^2}} - 1 \right)}{\alpha^2 \left(1 + \frac{2}{\alpha} + \frac{1}{\alpha^2} \right)} \\ &= \lim_{\alpha \rightarrow \infty} \frac{\alpha^2 \left(\sqrt{1 + \frac{2}{\alpha} + \frac{2}{\alpha^2}} - \frac{1}{\alpha} \right)}{\alpha^2 \left(1 + \frac{2}{\alpha} + \frac{1}{\alpha^2} \right)} \\ &= \lim_{\alpha \rightarrow \infty} \frac{\sqrt{1 + \frac{2}{\alpha} + \frac{2}{\alpha^2}} - \frac{1}{\alpha}}{1 + \frac{2}{\alpha} + \frac{1}{\alpha^2}} \\ &= \lim_{\alpha \rightarrow \infty} \frac{\sqrt{1 + \overset{0}{\frac{2}{\alpha}} + \overset{0}{\frac{2}{\alpha^2}}} - \overset{0}{\frac{1}{\alpha}}}{1 + \overset{0}{\frac{2}{\alpha}} + \overset{0}{\frac{1}{\alpha^2}}} \\ &= \frac{\sqrt{1}}{1} = 1 \end{aligned}$$

$$\boxed{\therefore \lim_{\alpha \rightarrow \infty} f(\alpha) = 1}$$

3.

$$\begin{aligned} \frac{\alpha \left(\sqrt{(\alpha + 1)^2 + 1} - 1 \right)}{(\alpha + 1)^2} &= \frac{1}{2} \\ \sqrt{(\alpha + 1)^2 + 1} &= \frac{(\alpha + 1)^2}{2\alpha} + 1 \\ (\alpha + 1)^2 + 1 &= \left(\frac{(\alpha + 1)^2}{2\alpha} + 1 \right)^2 \\ \alpha^2 + 2\alpha + 2 &= \frac{((\alpha + 1)^2 + 2\alpha)^2}{4\alpha^2} \\ 4\alpha^4 + 8\alpha^3 + 8\alpha^2 &= (\alpha^2 + 4\alpha + 1)^2 \\ 4\alpha^4 + 8\alpha^3 + 8\alpha^2 &= \alpha^4 + 8\alpha^3 + 18\alpha^2 + 8\alpha + 1 \\ 3\alpha^4 - 10\alpha^2 - 8\alpha - 1 &= 0 \end{aligned}$$

By the zero factor principle, $\alpha = -1$ is a potential solution:

$$3(-1)^4 - 10(-1)^2 - 8(-1) - 1 = 3 - 10 + 8 - 1 = 0$$

Knowing this, we can use long division to factor the polynomial.

$$\begin{array}{r} \overline{3\alpha^3 - 3\alpha^2 - 7\alpha - 1} \\ \alpha + 1 \overline{) 3\alpha^4 + 0\alpha^3 - 10\alpha^2 - 8\alpha - 1} \\ \underline{- 3\alpha^4 + 3\alpha^3} \\ \overline{-3\alpha^3 - 10\alpha^2} \\ \underline{- -3\alpha^3 - 3\alpha^2} \\ \phantom{\overline{-3\alpha^3 - 10\alpha^2}} -7\alpha^2 - 8\alpha \\ \phantom{\overline{-3\alpha^3 - 10\alpha^2}} \underline{- -7\alpha^2 - 7\alpha} \\ \phantom{\overline{-3\alpha^3 - 10\alpha^2}} -\alpha - 1 \\ \phantom{\overline{-3\alpha^3 - 10\alpha^2}} \underline{- -\alpha - 1} \\ \phantom{\overline{-3\alpha^3 - 10\alpha^2}} 0 \end{array}$$

By the same principle, $\alpha = -1$ is yet again another potential solution, thus $\alpha + 1$ is

another factor:

$$3(-1)^3 - 3(-1)^2 - 7(-1) - 1 = -3 - 3 + 7 - 1 = 0$$

If we use long division again,

$$(3\alpha^3 - 3\alpha^2 - 7\alpha - 1) \div (\alpha + 1) = 3\alpha^2 - 6\alpha - 1$$

Therefore, when we factor the polynomial we get:

$$3\alpha^4 - 10\alpha^2 - 8\alpha - 1 = (\alpha + 1)^2(3\alpha^2 - 6\alpha - 1)$$

If we use the the quadratic equation to solve the last factor, we get:

$$\begin{aligned}\alpha &= \frac{6 \pm \sqrt{(-6)^2 - 4(3)(-1)}}{2(3)} \\ &= \frac{6 \pm \sqrt{36 + 12}}{6} \\ &= \frac{6 \pm \sqrt{48}}{6} \\ &= \frac{6 \pm 4\sqrt{3}}{6} \\ &= \frac{3 \pm 2\sqrt{3}}{3}\end{aligned}$$

Therefore, our potential solutions are

$$\alpha = -1, \frac{3 \pm 2\sqrt{3}}{3}$$

However, $\alpha = -1$ or $\alpha = \frac{3 - 2\sqrt{3}}{3}$ are neither mathematically true nor physically possible. Therefore, to get exactly half the half-range, or $f(\alpha) = \frac{1}{2}$,

$$\boxed{\alpha = \frac{3 + 2\sqrt{3}}{3} \approx 2.2 \text{ times more massive than } M}$$